

# Explicit Formulas for Relaxed Disarrangement Densities from Structured Deformations

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**CF** a SD is a pair  $(g, G)$  with  $Dg = \nabla g \mathcal{L}^n + [g] \otimes \nu \mathcal{H}^{N-1}$ . **Approx. Theorem:** there exists  $f_n \in SBV$  such that  $f_n \xrightarrow{L^1} g$ ,  $\nabla f_n \xrightarrow{M} G$ .



# Energies arising from structured deformations

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Substituting  $|\cdot|$  with  $(\cdot)^\pm$  implies

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hence the question: does  $\mathcal{V}^{|\cdot|}$  have an associated disarrangement density? If so, how does it depend on  $(g, G)$ ?



# Relaxation *à la* Choksi-Fonseca - I

Given an energy density of the form

$$E(u) := \int_{\Omega} W(\nabla u) \, d\mathcal{L}^N + \int_{J(u) \cap \Omega} \psi([u], \nu) \, d\mathcal{H}^{N-1},$$





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$$I(g, G) = \int_{\Omega} H(\nabla g, G) \, d\mathcal{L}^N + \int_{J(g) \cap \Omega} h([g], \nu) \, d\mathcal{H}^{N-1}.$$



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The bulk density  $H$  and the surface density  $h$  are derived via a *cell formula*.



# Relaxation *à la* Choksi-Fonseca - II

## Theorem (Owen-Paroni, 2015)

Given  $\psi : \mathbb{R}^N \times \mathcal{S}^{N-1} \rightarrow [0, +\infty)$  such that

- there exists  $C > 0$ :  $0 \leq \psi(\xi, \nu) \leq C|\xi|$ ,
- for all  $t > 0$ ,  $\psi(t\xi, \nu) = t\psi(\xi, \nu)$ ,
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if we define, for any  $p > 1$ ,

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we have

$$I(\mathbf{g}, \mathbf{G}) = \int_{\Omega} H(\nabla \mathbf{g}, \mathbf{G}) \, d\mathcal{L}^N + \int_{J(\mathbf{g}) \cap \Omega} h([\mathbf{g}], \nu) \, d\mathcal{H}^{N-1}.$$

## Relaxation *à la* Choksi-Fonseca - III

### Theorem (Owen-Paroni, 2015)

The densities  $H$  and  $h$  are given by

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where

$$u_{\xi, \eta}(x) := \begin{cases} 0 & \text{if } -1/2 < x \cdot \eta < 0, \\ \xi & \text{if } 0 \leq x \cdot \eta < 1/2. \end{cases}$$

# Explicit formulas

## Theorem (Owen-Paroni (2015))

The choices  $\psi^{|\cdot|}(\xi, \nu) := |\xi \cdot \nu|$  and  $\psi^\pm(\xi, \nu) := (\xi \cdot \nu)^\pm$  yield

$$H^{|\cdot|}(A, B) = |\operatorname{tr}(A - B)|, \quad h^{|\cdot|}(\xi, \nu) = |\xi \cdot \nu| = \psi^{|\cdot|}(\xi, \nu),$$

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Therefore, the functional for the *minimal volume swept out by disarrangements* reads

$$\mathcal{V}^{|\cdot|}(g, G; \mathcal{P}) = \int_{\mathcal{P}} |\operatorname{tr}(\nabla g - G)| \, d\mathcal{L}^N + \int_{J(g) \cap \mathcal{P}} |[g] \cdot \nu| \, d\mathcal{H}^{N-1}.$$



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A crucial step in OP's proof was to show that

$$H^{|\cdot|}(A, B) \leq |\operatorname{tr}(A - B)|.$$



# Our contribution

Given  $A, B \in \mathbb{R}^{N \times N}$ , there holds

$$|\operatorname{tr}(A - B)| \leq \inf \left\{ \int_{J(u)} |[u] \cdot \nu| \, d\mathcal{H}^{N-1} : u \in SBV(Q; \mathbb{R}^N) \right. \\ \left. u(x) = Ax \text{ on } \partial Q, \nabla u \in L^p(Q), \int_Q \nabla u = B \right\}$$



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We prove that

$$\leq |\operatorname{tr}(A - B)|.$$





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- **Matias' Lemma (2007)**: Given  $\Omega \subset \mathbb{R}^N$  bounded open and  $M \in \mathbb{R}^{N \times N}$ , there exists  $C(N) > 0$  and  $u \in SBV(\Omega; \mathbb{R}^N)$ :



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Thank you for your attention!

